

# Conformal Bootstrap: a dream come true

Alessandro Vichi



13th March, 2015  
Bay Area Particle Theory Seminar

San Francisco State University

The study of asymptotic behaviors plays a central role in QFT, especially IR fixed points (universality)

- ▶ All examples where the IR behavior is known correspond to conformal invariant fixed points.
- ▶ In 4D, if perturbative, a fixed point is a CFT
- ▶ If non-perturbative, no formal proof but conformality largely accepted.
- ▶ In 2D, scale invariance implies conformal invariance

A large class of physically interesting IR fixed points are:

- ▶ non-supersymmetric
- ▶ non-perturbative
- ▶ small-N

No need to hunt for such a model... write the simplest (non-free) QFT:

$$\frac{1}{2}(\partial_\mu\phi)^2 + \frac{1}{2}\phi^2 + \frac{1}{4!}\phi^4, \quad \text{in } 2 \leq D < 4$$

How do we describe the properties of the Wilson-Fisher fixed point, say in 3D?

Is it stable? (namely, are there relevant operators singlet under global symmetries?)

## Goal of conformal bootstrap

*To develop a systematic and rigorous method to study the properties of conformal invariant fixed points*

- 1 What is a CFT?
- 2 Conformal bootstrap
- 3 An application: 3D ising model

## Conformal Algebra

In  $D$  dimensions :  $M_{\mu\nu}, P_\rho, D, K_\sigma \simeq SO(D|2)$

Irreducible representations of Conformal Algebra:

- ▶ infinite towers of states (or operators) with increasing, equally spaced, dimensions.
- ▶ Lower state is called **Primary**:

$$\mathcal{O}_{\Delta,\ell} : \begin{array}{cc} \Delta & \text{dimension} \\ \ell & \text{spin} \end{array}$$

- ▶ Other states, called **Descendants**, obtained applying  $P_\mu$
- ▶ representation totally characterized by **scaling dimension** and **spin** of the **primary**

Completeness of the Hilbert space of states  $\Leftrightarrow$  OPE:

$$\mathcal{O}_{\Delta_1}(x) \times \mathcal{O}_{\Delta_2}(y) = \frac{1}{|x-y|^{\Delta_1+\Delta_2}} \sum_{\mathcal{O}} C_{12\mathcal{O}} \underbrace{(C_{\mu_1 \dots \mu_\ell}(y) \mathcal{O}_{\Delta}^{\mu_1 \dots \mu_\ell}(y) + \text{descendants})}_{\text{fixed by conformal symmetry}}$$

$C_{12\mathcal{O}}$  are called **OPE coefficients**

## The power of conformal invariance

Two point function of primaries: completely fixed


$$\langle \mathcal{O}_i(x_1) \mathcal{O}_j(x_2) \rangle = \frac{\delta_{ij}}{x_{12}^{2\Delta_i}} \quad x_{12} \equiv |x_1 - x_2| \quad \Delta_i = [\mathcal{O}_i]$$

Three point function of primaries: fixed modulo a **constant**


$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \rangle \propto \begin{cases} C_{123} \underbrace{(\langle \mathcal{O}_3 \mathcal{O}_3 \rangle + \text{descendants})}_{\text{fixed by conformal symmetry}} & \text{if } \mathcal{O}_3 \in \mathcal{O}_1 \times \mathcal{O}_2 \\ 0 & \text{otherwise} \end{cases}$$

## Four point functions

Use OPE to reduce higher point functions to smaller ones

$$\langle \underbrace{\mathcal{O}(x_1)\mathcal{O}(x_2)} \underbrace{\mathcal{O}(x_3)\mathcal{O}(x_4)} \rangle \sim \sum_{\mathcal{O}} \text{diagram}$$


The diagram shows a four-point vertex where two external lines on the left meet at a point, and two external lines on the right meet at another point. These two vertices are connected by a horizontal red line segment. A red letter 'O' is placed above the red line segment, representing an operator in the s-channel expansion.

$$\langle \underbrace{\mathcal{O}(x_1)\mathcal{O}(x_2)\mathcal{O}(x_3)} \mathcal{O}(x_4) \rangle \sim \sum_{\mathcal{O}} \text{diagram}$$


The diagram shows a four-point vertex where two external lines on the left meet at a point, and two external lines on the right meet at another point. These two vertices are connected by a vertical red line segment. A red letter 'O' is placed to the left of the red line segment, representing an operator in the t-channel expansion.

If OPE is associative, the two expansion must give the same result!

## Definition of a CFT:

*A Conformal Field Theory is an infinite set of primary operators  $\mathcal{O}_{\Delta,\ell}$  and OPE coefficients  $C_{ijk}$  that satisfy crossing symmetry for all set of **four**-point functions.*



## Four point functions (more in details)

Recalling the OPE

$$\mathcal{O}(x_1) \times \mathcal{O}(x_2) = \sum_{\mathcal{O}'} \frac{C_{\mathcal{O}'}}{x_{12}^{2d-\Delta}} (\mathcal{O}'_{\Delta,\ell} + \text{descendants}) \quad d = [\mathcal{O}]$$

Then

$$\langle \underbrace{\mathcal{O}(x_1)\mathcal{O}(x_2)} \underbrace{\mathcal{O}(x_3)\mathcal{O}(x_4)} \rangle = \frac{u^{-d}}{(x_{13}^{2d} x_{24}^{2d})} \sum_{\mathcal{O}'_{\Delta,l}} C_{\mathcal{O}'}^2 \underbrace{\left( \langle \mathcal{O}'_{\Delta,\ell} \mathcal{O}'_{\Delta,\ell} \rangle + \text{descendants} \right)}_{\text{function of } u, v \text{ only by conformal symmetry}}$$

$$u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2} \quad v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$$

Conformal Blocks:

$$g_{\Delta,l}(u, v) \equiv \langle \mathcal{O}'_{\Delta,\ell} \mathcal{O}'_{\Delta,\ell} \rangle + \text{descendants}$$

They **sum up** the contribution of an **entire** representation

## The Bootstrap program

### Crossing Symmetry

$$\langle \underbrace{\mathcal{O}(x_1)\mathcal{O}(x_2)} \underbrace{\mathcal{O}(x_3)\mathcal{O}(x_4)} \rangle \quad \text{vs} \quad \langle \underbrace{\mathcal{O}(x_1)\mathcal{O}(x_2)\mathcal{O}(x_3)} \mathcal{O}(x_4) \rangle$$

They must produce the same result:

$$u^{-d} \left( 1 + \sum_{\Delta, l} C_{\Delta, l}^2 g_{\Delta, l}(u, v) \right) = v^{-d} \left( 1 + \sum_{\Delta, l} C_{\Delta, l}^2 g_{\Delta, l}(v, u) \right) \quad d = [\mathcal{O}]$$

Crossing symmetry  $\Rightarrow$  Sum Rule:

$$\sum_{\Delta, l} C_{\Delta, l}^2 \underbrace{\frac{v^d g_{\Delta, l}(u, v) - u^d g_{\Delta, l}(v, u)}{u^d - v^d}}_{F_{d, \Delta, l}} = 1$$

[Rattazzi, Rychkov, Tonni, AV]

- Breakthrough in the field in 2000: first computation of  $g_{\Delta, l}$  in  $D=2,4$
- At present  $g_{\Delta, l}$  are known numerically in any dimension
- Great efforts to extend to non scalar four point functions

## Geometric interpretation

$$\sum_{\Delta,\ell} C_{\Delta,l}^2 \begin{pmatrix} F_{d,\Delta,\ell} \\ \partial_u F_{d,\Delta,\ell} \\ \partial_v F_{d,\Delta,\ell} \\ \vdots \\ \partial_u^n \partial_v^m F_{d,\Delta,\ell} \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ \vdots \\ \vdots \end{pmatrix}$$

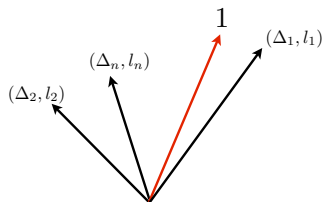
$F_{d,\Delta,\ell}$ : combinations of conformal blocks  
 $n + m \leq N_{\max}$



## Geometric interpretation

$$\sum_{\Delta, \ell} C_{\Delta, \ell}^2 \begin{pmatrix} F_{d, \Delta, \ell} \\ \partial_u F_{d, \Delta, \ell} \\ \partial_v F_{d, \Delta, \ell} \\ \vdots \\ \partial_u^n \partial_v^m F_{d, \Delta, \ell} \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ \vdots \\ \vdots \end{pmatrix}$$

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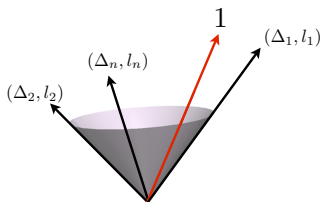


- All possible sums of vectors with positive coefficients define a cone

## Geometric interpretation

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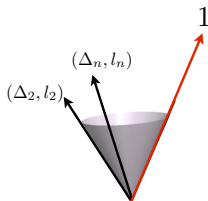


- All possible sums of vectors with positive coefficients define a cone
- Crossing symmetry satisfied  $\Leftrightarrow$  1 is inside the cone

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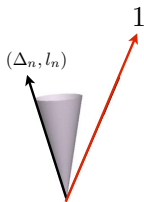


- All possible sums of vectors with positive coefficients define a cone
- Crossing symmetry satisfied  $\Leftrightarrow 1$  is inside the cone
- Restrictions on the spectrum make the cone narrower

## Geometric interpretation

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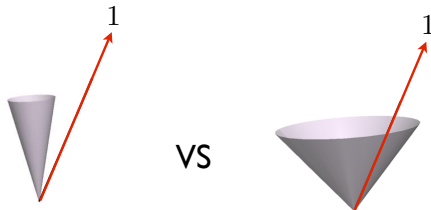
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- All possible sums of vectors with positive coefficients define a cone
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- Restrictions on the spectrum make the cone narrower
- A cone too narrow can't satisfy crossing symmetry: inconsistent spectrum

## Geometric interpretation

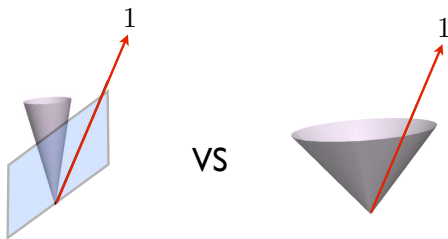
How can we distinguish feasible spectra from unfeasible ones?





## Geometric interpretation

How can we distinguish feasible spectra from unfeasible ones?



For unfeasible spectra it exists a plane separating the cone and the vector.

More formally...

Look for a Linear functional

$$\Lambda[F_{d,\Delta,\ell}] \equiv \sum_{n,m}^{N_{\max}} \lambda_{mn} \partial^n \partial^m F_{d,\Delta,\ell}$$

such that

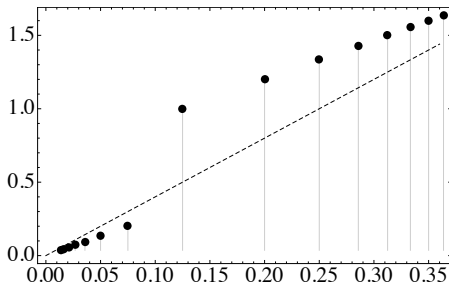
$$\Lambda[F_{d,\Delta,\ell}] > 0 \quad \text{and} \quad \Lambda[1] < 0$$

## 2D Example

Consider the OPE of scalar field in 2D CFT  $\phi$  with itself:

$$\phi \times \phi \sim 1 + \phi^2 + \text{higher dimensional operators,} \\ + \text{higher spin operators} \quad \Delta_\phi = [\phi], \quad \Delta_{\phi^2} = [\phi^2]$$

What values of  $(\Delta_\phi, \Delta_{\phi^2})$  are consistent with crossing symmetry?  
(black points are minimal models, exactly known CFT's)

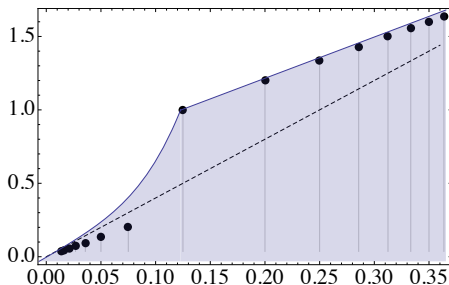


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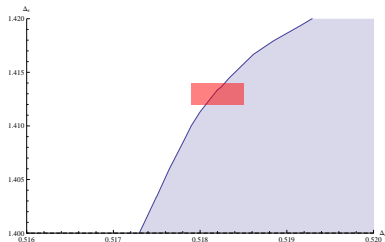
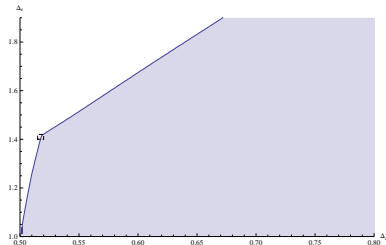
[Rychkov, AV]

## 3D Ising Model

Some notation:

$$\sigma \times \sigma \sim 1 + \epsilon + \epsilon' + \dots$$

Allowed regions in  $\Delta_\sigma, \Delta_\epsilon$  plane ?



[El-Showk,Paulos,Poland,Rychkov,Simmons-Duffin, AV]

Already excluding part of  $\epsilon$ -expansion prediction

## Going beyond: multiple correlators

So far we used a single four point function :  $\langle \sigma \sigma \sigma \sigma \rangle$ .

Let us include additional correlators:  $\langle \epsilon \epsilon \epsilon \epsilon \rangle$ ,  $\langle \sigma \epsilon \sigma \epsilon \rangle$ .

Must consider all mixed OPE's :

$$\begin{aligned} \sigma \times \sigma &\sim 1 + \epsilon + \epsilon' + \dots \quad \mathbb{Z}_2 - \text{even} \\ \sigma \times \epsilon &\sim \sigma + \sigma' + \dots \quad \mathbb{Z}_2 - \text{odd} \\ \epsilon \times \epsilon &\sim 1 + \epsilon + \epsilon' + \dots \quad \mathbb{Z}_2 - \text{even} \end{aligned}$$

$$\begin{aligned} \langle \sigma(x_1) \epsilon(x_2) \sigma(x_3) \epsilon(x_4) \rangle &\sim \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\sigma\epsilon\mathcal{O}}^2 \tilde{g}_{\Delta,\ell}(u,v) \\ \langle \sigma(x_1) \epsilon(x_2) \sigma(x_3) \epsilon(x_4) \rangle &\sim \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\sigma\sigma\mathcal{O}} \lambda_{\epsilon\epsilon\mathcal{O}} g_{\Delta,\ell}(u,v) \end{aligned}$$

Second expansion is not a sum with positive coefficients: geometrical argument can't go through, but it can be generalized.

*Study region allowed by multi-correlators crossing symmetry under the unique assumption that  $\sigma$  and  $\epsilon$  are the only two relevant scalar operators in theory.*

## 3D Ising Model: multiple correlators

Some notation:

$$\sigma \times \sigma \sim 1 + \epsilon + \epsilon' + \dots \quad \mathbb{Z}_2 - \text{even}$$

$$\sigma \times \epsilon \sim \sigma + \sigma' + \dots \quad \mathbb{Z}_2 - \text{odd}$$

$$\epsilon \times \epsilon \sim 1 + \epsilon + \epsilon' + \dots \quad \mathbb{Z}_2 - \text{even}$$

Use  $\langle \sigma\sigma\sigma\sigma \rangle, \langle \sigma\sigma\epsilon\epsilon \rangle, \langle \epsilon\epsilon\epsilon\epsilon \rangle$

## 3D Ising Model: multiple correlators

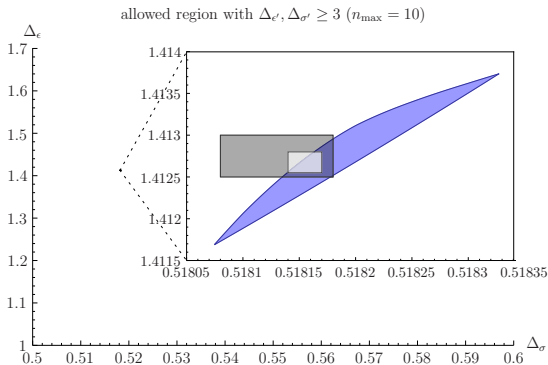
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[Kos,Poland,Simmons-Duffin]



### 3D Ising Model: multiple correlators

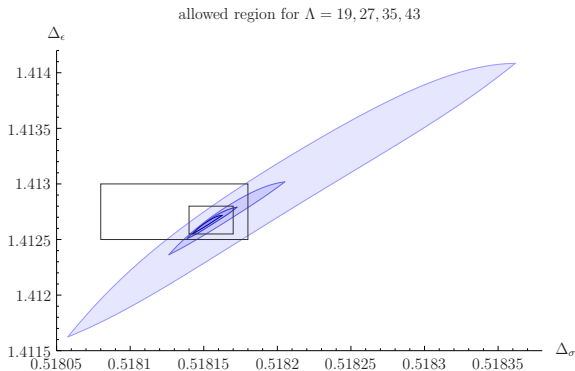
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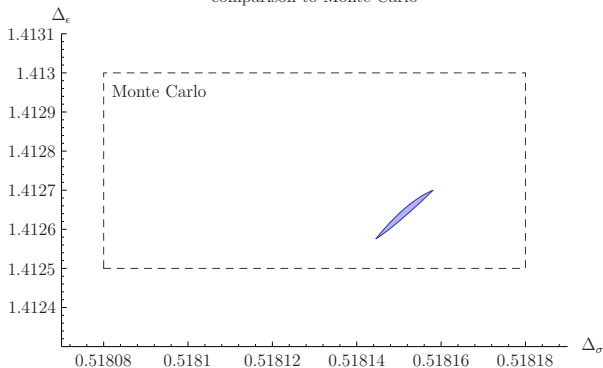
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Use  $\langle \sigma\sigma\sigma\sigma \rangle, \langle \sigma\sigma\epsilon\epsilon \rangle, \langle \epsilon\epsilon\epsilon\epsilon \rangle$

comparison to Monte Carlo



[Simmons-Duffin]

$\Delta_\sigma \in [0.518145, 0.518157]$

## Summary

spin & $\mathbb{Z}_2$	name	$\Delta$	OPE coefficient
$\ell = 0, \mathbb{Z}_2 = -$	$\sigma$	0.518145(6)	
$\ell = 0, \mathbb{Z}_2 = +$	$\epsilon$	1.41264(6)	$f_{\sigma\sigma\epsilon}^2 = 1.10636(9)$
	$\epsilon'$	3.8303(18)	$f_{\sigma\sigma\epsilon'}^2 = 0.002810(6)$
$\ell = 2, \mathbb{Z}_2 = +$	$T$	3	$c_T/c_T^{\text{free}} = 0.946534(11)$
	$T'$	5.500(15)	$f_{\sigma\sigma T'}^2 = 2.97(2) \times 10^{-4}$

- The 3D Ising model is a CFT with only two relevant operators:  $\sigma$  and  $\epsilon$
- The 3D Ising model lies on the boundary of the region allowed by single correlator crossing symmetry
- Operator dimensions give the most precise determination of  $\nu, \eta, \omega$  critical exponents to date

$$\Delta_\sigma = 1/2 + \eta/2 \quad \Delta_\epsilon = 3 - 1/\nu \quad \Delta_{\epsilon'} = 3 + \omega \quad \Delta_{\epsilon''} = 3 + \omega_2 \quad \Delta_{\epsilon'''} = 3 + \omega_3$$

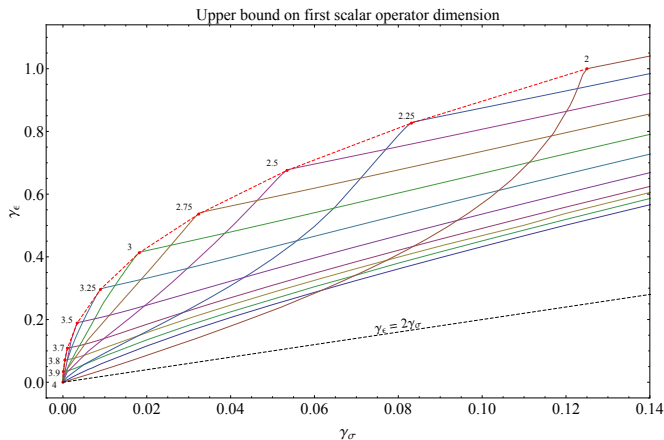
- First precise estimate of OPE coefficients and central charge
- Additional operators and coefficients (with larger errorbars) can be extracted
- What next? multiple correlators analysis can pinpoint the location of  $O(N)$ -models [F. Kos, D. Poland, D. Simmons-Duffin, AV, in progress]
- Study correlation functions containing conserved currents [ AV, in progress ; M Costa & al, in progress]

# BACKUP SLIDES

## A proliferation of kinks

Compare bounds on the **anomalous dimensions** for various space-time dimensions  $D$ :

$$\gamma_\sigma = \Delta_\sigma - \frac{(D-2)}{2} \quad \gamma_\epsilon = \Delta_\epsilon - (D-2)$$

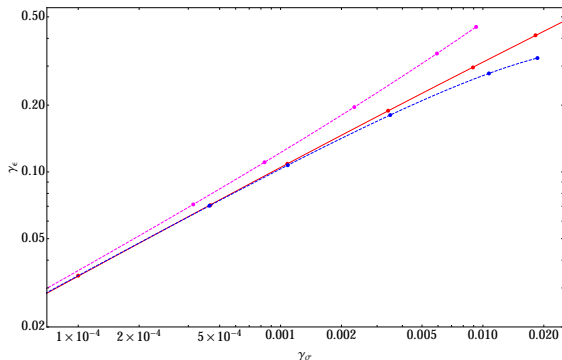


**Epsilon Expansion:**  $D = 4 - \varepsilon$ 

$$\gamma_\sigma = \frac{(N+2)\varepsilon^2}{4(N+8)^2} + O(\varepsilon^3)$$

$$\gamma_\epsilon = \frac{(N+2)\varepsilon}{N+8} + \frac{(N+2)(13N+44)\varepsilon^2}{2(N+8)^3} + O(\varepsilon^3)$$

Comparison with epsilon-expansion at 2–3 loops

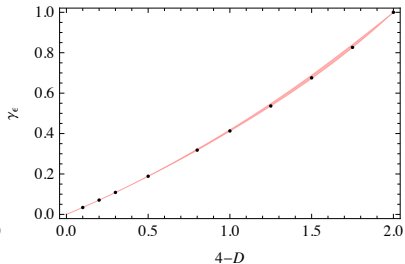
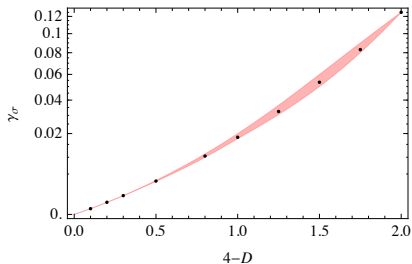


► Kinks from previous slide

►  $O(\varepsilon^2)$

►  $O(\varepsilon^3)$

## Epsilon Expansion: $D = 4 - \varepsilon$



- Our prediction (points)
- Borel resummed series: central values and errors (bands)

[Guillou, Zinn-Justin]

## Multiple correlators

When using multiple correlators the search for linear functionals must be modified to accommodate non squared OPE coefficients:

Single correlator:

$$\sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\sigma\sigma\mathcal{O}}^2 F_{\Delta\sigma,\Delta,\ell} = 1$$

Look for functional

$$\Lambda[F_{\Delta\sigma,\Delta,\ell}] \equiv \sum_{n,m}^{N_{\max}} \lambda_{mn} \partial^n \partial^m F_{\Delta\sigma,\Delta,\ell}$$

such that:

$$\Lambda[F_{\Delta\sigma,\Delta,\ell}] > 0 \quad \text{and} \quad \Lambda[1] < 0$$

Multi correlators:

$$\sum_{\mathcal{O}_{\Delta,\ell}} \vec{\lambda}_{\mathcal{O}}^T M_{\Delta,\ell} \vec{\lambda}_{\mathcal{O}} + \sum_{\mathcal{O}'_{\Delta,\ell}} \lambda_{\sigma\epsilon\mathcal{O}'}^2 \tilde{F}_{\Delta,\ell} = 0$$

$$M_{\Delta,\ell} = \begin{pmatrix} 0 & \frac{1}{2} F_{\Delta\sigma,\Delta,\ell} \\ \frac{1}{2} F_{\Delta\epsilon,\Delta,\ell} & 0 \end{pmatrix}, \quad \vec{\lambda}_{\mathcal{O}} = \begin{pmatrix} \lambda_{\sigma\sigma\mathcal{O}} \\ \lambda_{\epsilon\epsilon\mathcal{O}} \end{pmatrix}$$

Look for a functional acting on matrices

$$\Lambda[M_{\Delta,\ell}] \equiv \sum_{n,m}^{N_{\max}} \lambda_{mn} \partial^n \partial^m M_{\Delta,\ell}$$

such that (semidefinite positiveness condition)

$$\Lambda[M_{\Delta,\ell}] \succeq 0 \quad \text{and} \quad \Lambda[M_{0,0}] < 0$$



## Solving Ising 3D: boundary spectrum

- ▶ On the boundary of the allowed region the solution to crossing is unique: the whole spectrum and OPE coefficients can be reconstructed.
- ▶ Assuming to leave on the upper boundary of the allowed island (note that increasing the numerical power it is approximatively stable)

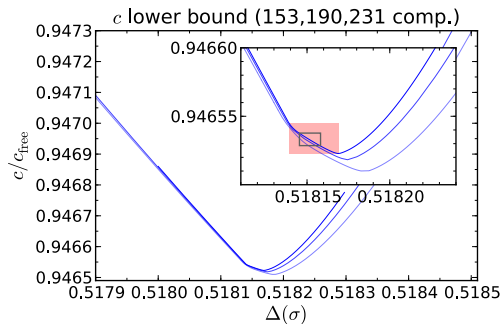
Recall OPE:  $\sigma \times \sigma \sim 1 + \epsilon + \epsilon' + \dots (\ell=0)$   
 $+ T_{\mu\nu} + T'_{\mu\nu} + \dots (\ell=2)$   
 $+ \dots (\ell > 2)$

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 $+ \dots$  ( $\ell > 2$ )

Prediction for central charge:



$$c/c_{\text{free}} \in [0.946528, 0.946538]$$

red rectangle : assuming Ising 3D has minimal central charge

gray rectangle: multiple correlators

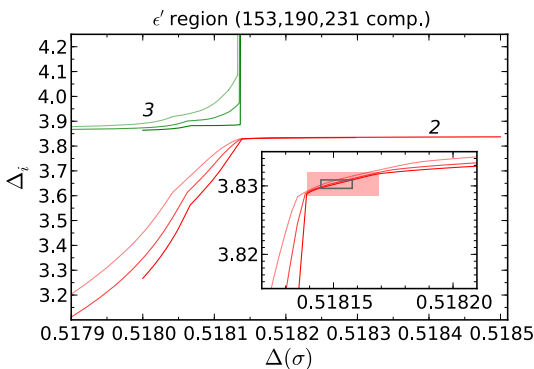
[El-Showk,Paulos,Poland,Rychkov,Simmons-Duffin, AV]

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 $+ T_{\mu\nu} + T'_{\mu\nu} + \dots$  ( $\ell=2$ )  
 $+ \dots$  ( $\ell > 2$ )

Prediction for  $\epsilon'$ :



$$\Delta_{\epsilon'} \in [3.829, 3.831]$$

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gray rectangle: multiple correlators

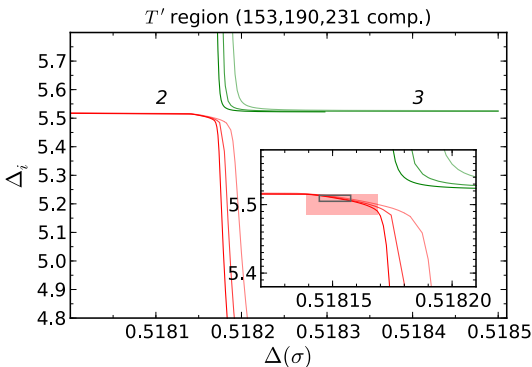
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## Solving Ising 3D: boundary spectrum

- ▶ On the boundary of the allowed region the solution to crossing is unique: the whole spectrum and OPE coefficients can be reconstructed.
- ▶ Assuming to leave on the upper boundary of the allowed island (note that increasing the numerical power it is approximatively stable)

Recall OPE:  $\sigma \times \sigma \sim 1 + \epsilon + \epsilon' + \dots$  ( $\ell=0$ )  
 $+ T_{\mu\nu} + T'_{\mu\nu} + \dots$  ( $\ell=2$ )  
 $+ \dots$  ( $\ell > 2$ )

Prediction for  $T'$ :



$$\Delta_{T'} \in [5.505, 5.515]$$

red rectangle : assuming Ising 3D has minimal central charge  
 gray rectangle: multiple correlators

[El-Showk,Paulos,Poland,Rychkov,Simmons-Duffin, AV]